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Introduction to the Calculus of Variations: Lecture 4

Swarnendu Sil

Department of Mathematics Indian Institute of Science

Spring Semester 2021

Outline

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As we mentioned, that H is a conserved quantity if H does not depend explicitly on t, is not a coincidence, but is just an example of a profound general fact.

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As we mentioned, that H is a conserved quantity if H does not depend explicitly on t, is not a coincidence, but is just an example of a profound general fact.

Symmetries of \mathcal{L} or $H \Leftrightarrow$ first integral.

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As we mentioned, that H is a conserved quantity if H does not depend explicitly on t, is not a coincidence, but is just an example of a profound general fact.

Symmetries of \mathcal{L} or $H \Leftrightarrow$ first integral.

Any symmetry of the Lagrangian (and thus also of the Hamiltonian and vice versa) corresponds to a first integral, i.e. a conserved quantity.

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This is the Noether's theorem, named after its discoverer, the brilliant Emmy Noether.

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I spent half a slide for a reason, to remind you something that I sincerely hope should not need reminding – that mathematicians do not have to be men and genius does not have a gender.

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Now let us show that if the Hamiltonian (and thus the Lagrangian too) does not depend explicitly on t, then there is a symmetry.

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Now let us show that if the Hamiltonian (and thus the Lagrangian too) does not depend explicitly on t, then there is a symmetry. This symmetry is called the time translation symmetry or in other words, invariance under time translations.

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Now let us show that if the Hamiltonian (and thus the Lagrangian too) does not depend explicitly on t, then there is a symmetry. This symmetry is called the time translation symmetry or in other words, invariance under time translations. Consider the

one-parameter family of diffeomorphisms

$$\phi_{\tau}(t) = t + \tau$$
 for $\tau \in \mathbb{R}$.

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$$\phi_{\tau}(t) = t + \tau$$
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Thus, the curve $u : [a, b] \to \mathbb{R}^N$ is transformed to $u \circ \phi_{\tau}^{-1} : [a + \tau, b + \tau] \to \mathbb{R}^N$.

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Thus, the curve $u : [a, b] \to \mathbb{R}^N$ is transformed to $u \circ \phi_{\tau}^{-1} : [a + \tau, b + \tau] \to \mathbb{R}^N$. Now, since the Hamiltonian does not depend explicitly on t, the same must be true for the Lagrangian density f Introduction to the Calculus of Variations

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Now let us show that if the Hamiltonian (and thus the Lagrangian too) does not depend explicitly on t, then there is a symmetry. This symmetry is called the time translation symmetry or in other words, invariance under time translations. Consider the one-parameter family of diffeomorphisms

$$\phi_{ au}(t) = t + au$$
 for $au \in \mathbb{R}$.

Thus, the curve $u : [a, b] \to \mathbb{R}^N$ is transformed to $u \circ \phi_{\tau}^{-1} : [a + \tau, b + \tau] \to \mathbb{R}^N$. Now, since the Hamiltonian does not depend explicitly on t, the same must be true for the Lagrangian density f and thus we have, for any $\theta \in (a, b]$,

$$\int_{a}^{\theta} f\left(u\left(t\right), \dot{u}\left(t\right)\right) \, \mathrm{d}t = \int_{a+\tau}^{\theta+\tau} f\left(u \circ \phi_{\tau}^{-1}\left(s\right), \frac{d}{ds}\left[u \circ \phi_{\tau}^{-1}\right]\left(s\right)\right) \, \mathrm{d}s.$$

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Thus, the curve $u : [a, b] \to \mathbb{R}^N$ is transformed to $u \circ \phi_{\tau}^{-1} : [a + \tau, b + \tau] \to \mathbb{R}^N$. Now, since the Hamiltonian does not depend explicitly on t, the same must be true for the Lagrangian density f and thus we have, for any $\theta \in (a, b]$,

$$\int_{a}^{\theta} f\left(u\left(t\right), \dot{u}\left(t\right)\right) \, \mathrm{d}t = \int_{a+\tau}^{\theta+\tau} f\left(u \circ \phi_{\tau}^{-1}\left(s\right), \frac{d}{ds}\left[u \circ \phi_{\tau}^{-1}\right]\left(s\right)\right) \, \mathrm{d}s.$$

Thus, we have the invariance

 $I_{ heta}[u] = I_{ heta}[u \circ \phi_{ au}^{-1}]$ for any $heta \in (a, b].$

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Definition (Invariance)

Let $\phi_s : [a, b] \times \mathbb{R}^N \to \mathbb{R} \times \mathbb{R}^N$, $s \in \mathbb{R}$ be a smoothly varying one-parameter family of diffeomorphisms,

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$$\phi_{s}\left(t,z\right)=\left(\phi_{s}^{0}\left(t\right),\bar{\phi}_{s}\left(z\right)\right) \ \, \text{for every} \ t\in\left[a,b\right] \text{ and for every} \ z\in\mathbb{R}^{N},$$

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$$\phi_{s}\left(t,z\right)=\left(\phi_{s}^{0}\left(t\right),\bar{\phi}_{s}\left(z\right)\right) \ \, \text{for every} \ t\in\left[a,b\right] \text{ and for every} \ z\in\mathbb{R}^{N},$$

such that

$$\phi_0(t,z) = (t,z)$$
 for every $t \in [a,b]$ and for every $z \in \mathbb{R}^N$

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$$\phi_{s}\left(t,z\right)=\left(\phi_{s}^{0}\left(t\right),\bar{\phi}_{s}\left(z\right)\right) \ \, \text{for every} \ t\in\left[a,b\right] \text{ and for every} \ z\in\mathbb{R}^{N},$$

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A Lagrangian is invariant under the action of the diffeomorphisms if it satisfies

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$$\phi_{s}\left(t,z
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ight),ar{\phi}_{s}\left(z
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A Lagrangian is invariant under the action of the diffeomorphisms if it satisfies

$$\begin{split} \int_{\phi_s^0(a)}^{\phi_s^0(\theta)} f\left(t_s, \left[\bar{\phi}_s \circ u \circ \left(\phi_s^0\right)^{-1}\right](t_s), \frac{d}{dt_s}\left[\bar{\phi}_s \circ u \circ \left(\phi_s^0\right)^{-1}\right](t_s)\right) \, \mathrm{d}t_s \\ &= \int_a^\theta f\left(t, u\left(t\right), \dot{u}\left(t\right)\right) \, \mathrm{d}t \qquad \text{for every } \theta \in (a, b], \end{split}$$

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$$\phi_{s}\left(t,z
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for every $s \in \mathbb{R}$,

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$$\phi_{s}\left(t,z\right)=\left(\phi_{s}^{0}\left(t\right),\bar{\phi}_{s}\left(z\right)\right) \ \, \text{for every} \ t\in\left[a,b\right] \text{ and for every} \ z\in\mathbb{R}^{N},$$

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A Lagrangian is invariant under the action of the diffeomorphisms if it satisfies

$$\begin{split} \int_{\phi_s^0(\theta)}^{\phi_s^0(\theta)} f\left(t_s, \left[\bar{\phi}_s \circ u \circ \left(\phi_s^0\right)^{-1}\right](t_s), \frac{d}{dt_s}\left[\bar{\phi}_s \circ u \circ \left(\phi_s^0\right)^{-1}\right](t_s)\right) \, \mathrm{d}t_s \\ &= \int_a^\theta f\left(t, u\left(t\right), \dot{u}\left(t\right)\right) \, \mathrm{d}t \qquad \text{for every } \theta \in (a, b], \end{split}$$

for every $s \in \mathbb{R}$, where $t_s = \phi_s^0(t)$.

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Theorem (Noether's theorem) Let $f \in C^2([a, b] \times \mathbb{R}^N \times \mathbb{R}^N)$.

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$$I[u] = \int_{a}^{b} f(t, u(t), \dot{u}(t)) dt$$

is invariant under the action of the diffeomorphisms $\{\phi_s\}$ as above.

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$$I[u] = \int_{a}^{b} f(t, u(t), \dot{u}(t)) dt$$

is invariant under the action of the diffeomorphisms $\{\phi_{\rm s}\}$ as above. Then the following expression

$$\left\langle f_{\xi}\left(t, u\left(t\right), \dot{u}\left(t\right)\right), \frac{d}{ds}\left[\bar{\phi}_{s}\left(u(t)\right)\right] \bigg|_{s=0} \right\rangle$$

+ $\left[f\left(t, u\left(t\right), \dot{u}\left(t\right)\right) - \left\langle f_{\xi}\left(t, u\left(t\right), \dot{u}\left(t\right)\right), \dot{u}\left(t\right)\right\rangle\right] \frac{d}{ds}\left[\phi_{s}^{0}\left(t\right)\right] \bigg|_{s=0}$

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$$I[u] = \int_{a}^{b} f(t, u(t), \dot{u}(t)) dt$$

is invariant under the action of the diffeomorphisms $\{\phi_{\rm s}\}$ as above. Then the following expression

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is constant along any solution u(t) of the EL equations for I,

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$$I[u] = \int_{a}^{b} f(t, u(t), \dot{u}(t)) dt$$

is invariant under the action of the diffeomorphisms $\{\phi_s\}$ as above. Then the following expression

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is **constant** along any solution u(t) of the EL equations for I, i.e. defines a first integral.

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First we prove a special case where the t variable is unchanged,

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First we prove a special case where the *t* variable is unchanged, i.e. $\phi_s^0(t) = t$ for every $t \in [a, b]$ for every $s \in \mathbb{R}$.

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$$\int_{a}^{\theta} f\left(t, \left[\bar{\phi}_{s} \circ u\right](t), \frac{d}{dt}\left[\bar{\phi}_{s} \circ u\right](t)\right) dt$$
$$= \int_{a}^{\theta} f\left(t, u\left(t\right), \dot{u}\left(t\right)\right) dt$$

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$$= \int_{a}^{\theta} f\left(t, u\left(t\right), \dot{u}\left(t\right)\right) dt$$

for every $s \in \mathbb{R}$, for every $\theta \in (a, b]$ and for any $u \in C^2([a, b]; \mathbb{R}^N)$.

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for every $s \in \mathbb{R}$, for every $\theta \in (a, b]$ and for any $u \in C^2([a, b]; \mathbb{R}^N)$. Note that a consequence of invariance, which we would not use in the proof, is that if u satisfies the EL equations, so does $\overline{\phi}_s \circ u$ for every $s \in \mathbb{R}$.

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for every $s \in \mathbb{R}$, for every $\theta \in (a, b]$ and for any $u \in C^2([a, b]; \mathbb{R}^N)$. Note that a consequence of invariance, which we would not use in the proof, is that if u satisfies the EL equations, so does $\overline{\phi}_s \circ u$ for every $s \in \mathbb{R}$. Thus, we have,

$$\frac{d}{dt} \left[f_{\xi} \left(t, \left[\bar{\phi}_{s} \circ u \right] \left(t \right), \frac{d}{dt} \left[\bar{\phi}_{s} \circ u \right] \left(t \right) \right) \right] \\
= f_{u} \left(t, \left[\bar{\phi}_{s} \circ u \right] \left(t \right), \frac{d}{dt} \left[\bar{\phi}_{s} \circ u \right] \left(t \right) \right)$$

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$$= \int_{a}^{\theta} f\left(t, u\left(t\right), \dot{u}\left(t\right)\right) dt$$

for every $s \in \mathbb{R}$, for every $\theta \in (a, b]$ and for any $u \in C^2([a, b]; \mathbb{R}^N)$. Note that a consequence of invariance, which we would not use in the proof, is that if u satisfies the EL equations, so does $\overline{\phi}_s \circ u$ for every $s \in \mathbb{R}$. Thus, we have,

$$\frac{d}{dt} \left[f_{\xi} \left(t, \left[\bar{\phi}_{s} \circ u \right] \left(t \right), \frac{d}{dt} \left[\bar{\phi}_{s} \circ u \right] \left(t \right) \right) \right] \\
= f_{u} \left(t, \left[\bar{\phi}_{s} \circ u \right] \left(t \right), \frac{d}{dt} \left[\bar{\phi}_{s} \circ u \right] \left(t \right) \right)$$

for every $s \in \mathbb{R}$ and for every $t \in (a, \theta)$. We would not use this.

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Fix $\theta \in (a, b]$.

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Examples
Fix $\theta \in (a, b]$. Now differentiating w.r.t s and using the fact that $\bar{\phi}_0$ is identity on \mathbb{R}^N ,

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$$0 = \left. \frac{d}{ds} \left(\int_{a}^{\theta} f\left(t, \left[\bar{\phi}_{s} \circ u\right](t), \frac{d}{dt} \left[\bar{\phi}_{s} \circ u\right](t) \right) \, \mathrm{d}t \right) \right|_{s=0}$$

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Fix $\theta \in (a, b]$. Now differentiating w.r.t s and using the fact that $\bar{\phi}_0$ is identity on \mathbb{R}^N , we deduce, by invariance,

$$0 = \left. \frac{d}{ds} \left(\int_{a}^{\theta} f\left(t, \left[\bar{\phi}_{s} \circ u\right](t), \frac{d}{dt} \left[\bar{\phi}_{s} \circ u\right](t) \right) dt \right) \right|_{s=0} \\ = \left. \int_{a}^{\theta} \left\langle f_{u}\left(t, u\left(t\right), \dot{u}\left(t\right)\right); \frac{d}{ds} \left[\bar{\phi}_{s} \circ u\right] \right|_{s=0} \left(t\right) \right\rangle dt$$

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$$= \int_{a}^{\theta} \left\langle f_{u}\left(t, u\left(t\right), \dot{u}\left(t\right)\right); \frac{d}{ds} \left[\bar{\phi}_{s} \circ u\right] \bigg|_{s=0}\left(t\right) \right\rangle dt$$
$$+ \int_{a}^{\theta} \left\langle f_{\xi}\left(t, u\left(t\right), \dot{u}\left(t\right)\right); \frac{d}{dt} \left[\frac{d}{ds} \left[\bar{\phi}_{s} \circ u\right] \bigg|_{s=0}\left(t\right)\right] \right\rangle dt$$

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$$0 = \frac{d}{ds} \left(\int_{a}^{\theta} f\left(t, \left[\bar{\phi}_{s} \circ u\right](t), \frac{d}{dt} \left[\bar{\phi}_{s} \circ u\right](t)\right) dt \right) \Big|_{s=0}$$
$$= \int_{a}^{\theta} \left\langle f_{u}\left(t, u\left(t\right), \dot{u}\left(t\right)\right); \frac{d}{ds} \left[\bar{\phi}_{s} \circ u\right] \Big|_{s=0}\left(t\right) \right\rangle dt$$
$$+ \int_{a}^{\theta} \left\langle f_{\xi}\left(t, u\left(t\right), \dot{u}\left(t\right)\right); \frac{d}{dt} \left[\frac{d}{ds} \left[\bar{\phi}_{s} \circ u\right] \Big|_{s=0}\left(t\right)\right] \right\rangle dt$$
$$= \int_{a}^{\theta} \left\langle \frac{d}{dt} \left[f_{\xi}\left(t, u\left(t\right), \dot{u}\left(t\right)\right)\right]; \frac{d}{ds} \left[\bar{\phi}_{s} \circ u\right] \Big|_{s=0}\left(t\right) \right\rangle dt$$

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$$0 = \frac{d}{ds} \left(\int_{a}^{\theta} f\left(t, \left[\bar{\phi}_{s} \circ u\right](t), \frac{d}{dt} \left[\bar{\phi}_{s} \circ u\right](t)\right) dt \right) \Big|_{s=0}$$
$$= \int_{a}^{\theta} \left\langle f_{u}\left(t, u\left(t\right), \dot{u}\left(t\right)\right); \frac{d}{ds} \left[\bar{\phi}_{s} \circ u\right] \Big|_{s=0}\left(t\right) \right\rangle dt$$
$$+ \int_{a}^{\theta} \left\langle f_{\xi}\left(t, u\left(t\right), \dot{u}\left(t\right)\right); \frac{d}{dt} \left[\frac{d}{ds} \left[\bar{\phi}_{s} \circ u\right] \Big|_{s=0}\left(t\right)\right] \right\rangle dt$$
$$= \int_{a}^{\theta} \left\langle \frac{d}{dt} \left[f_{\xi}\left(t, u\left(t\right), \dot{u}\left(t\right)\right)\right]; \frac{d}{ds} \left[\bar{\phi}_{s} \circ u\right] \Big|_{s=0}\left(t\right) \right\rangle dt$$

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$$\begin{split} 0 &= \left. \frac{d}{ds} \left(\int_{a}^{\theta} f\left(t, \left[\bar{\phi}_{s} \circ u \right](t), \frac{d}{dt} \left[\bar{\phi}_{s} \circ u \right](t) \right) \, \mathrm{d}t \right) \right|_{s=0} \\ &= \int_{a}^{\theta} \left\langle f_{u}\left(t, u\left(t \right), \dot{u}\left(t \right) \right); \frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \right|_{s=0}\left(t \right) \right\rangle \, \mathrm{d}t \\ &+ \int_{a}^{\theta} \left\langle f_{\xi}\left(t, u\left(t \right), \dot{u}\left(t \right) \right); \frac{d}{dt} \left[\left. \frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \right|_{s=0}\left(t \right) \right] \right\rangle \, \mathrm{d}t \\ &= \int_{a}^{\theta} \left\langle \frac{d}{dt} \left[f_{\xi}\left(t, u\left(t \right), \dot{u}\left(t \right) \right) \right]; \frac{d}{dt} \left[\left. \frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \right|_{s=0}\left(t \right) \right\rangle \, \mathrm{d}t \\ &+ \int_{a}^{\theta} \left\langle f_{\xi}\left(t, u\left(t \right), \dot{u}\left(t \right) \right) \right]; \frac{d}{dt} \left[\left. \frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \right|_{s=0}\left(t \right) \right\rangle \, \mathrm{d}t \end{split}$$

In the last line we substituted for f_u using the EL equations.

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$$0 = \int_{a}^{\theta} \left\langle \frac{d}{dt} \left[f_{\xi} \left(t, u\left(t \right), \dot{u}\left(t \right) \right) \right]; \frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \bigg|_{s=0} (t) \right\rangle \, \mathrm{d}t \\ + \int_{a}^{\theta} \left\langle f_{\xi} \left(t, u\left(t \right), \dot{u}\left(t \right) \right); \frac{d}{dt} \left[\frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \bigg|_{s=0} (t) \right] \right\rangle \, \mathrm{d}t$$

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$$0 = \int_{a}^{\theta} \left\langle \frac{d}{dt} \left[f_{\xi} \left(t, u\left(t \right), \dot{u}\left(t \right) \right) \right]; \frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \Big|_{s=0} \left(t \right) \right\rangle \, \mathrm{d}t \\ + \int_{a}^{\theta} \left\langle f_{\xi} \left(t, u\left(t \right), \dot{u}\left(t \right) \right); \frac{d}{dt} \left[\left[\frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \right|_{s=0} \left(t \right) \right] \right\rangle \, \mathrm{d}t \\ = \int_{a}^{\theta} \frac{d}{dt} \left[\left\langle f_{\xi} \left(t, u\left(t \right), \dot{u}\left(t \right) \right); \frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \right|_{s=0} \left(t \right) \right] \right\rangle \, \mathrm{d}t$$

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$$0 = \int_{a}^{\theta} \left\langle \frac{d}{dt} \left[f_{\xi} \left(t, u\left(t \right), \dot{u}\left(t \right) \right) \right]; \frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \Big|_{s=0} \left(t \right) \right\rangle \, \mathrm{d}t \\ + \int_{a}^{\theta} \left\langle f_{\xi} \left(t, u\left(t \right), \dot{u}\left(t \right) \right); \frac{d}{dt} \left[\left[\frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \right]_{s=0} \left(t \right) \right] \right\rangle \, \mathrm{d}t \\ = \int_{a}^{\theta} \frac{d}{dt} \left[\left\langle f_{\xi} \left(t, u\left(t \right), \dot{u}\left(t \right) \right); \frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \right|_{s=0} \left(t \right) \right] \right\rangle \, \mathrm{d}t$$

Thus, we have

$$\left\langle f_{\xi}\left(\theta, u\left(\theta\right), \dot{u}\left(\theta\right)\right); \frac{d}{ds} \left[\bar{\phi}_{s} \circ u\right] \Big|_{s=0}\left(\theta\right) \right\rangle$$
$$= \left\langle f_{\xi}\left(a, u\left(a\right), \dot{u}\left(a\right)\right); \frac{d}{ds} \left[\bar{\phi}_{s} \circ u\right] \Big|_{s=0}\left(a\right) \right\rangle.$$

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$$0 = \int_{a}^{\theta} \left\langle \frac{d}{dt} \left[f_{\xi} \left(t, u\left(t \right), \dot{u}\left(t \right) \right) \right]; \frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \Big|_{s=0} \left(t \right) \right\rangle \, \mathrm{d}t \\ + \int_{a}^{\theta} \left\langle f_{\xi} \left(t, u\left(t \right), \dot{u}\left(t \right) \right); \frac{d}{dt} \left[\left[\frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \right]_{s=0} \left(t \right) \right] \right\rangle \, \mathrm{d}t \\ = \int_{a}^{\theta} \frac{d}{dt} \left[\left\langle f_{\xi} \left(t, u\left(t \right), \dot{u}\left(t \right) \right); \frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \right|_{s=0} \left(t \right) \right] \right\rangle \, \mathrm{d}t$$

Thus, we have

$$\left\langle f_{\xi}\left(\theta, u\left(\theta\right), \dot{u}\left(\theta\right)\right); \frac{d}{ds} \left[\bar{\phi}_{s} \circ u\right] \Big|_{s=0}\left(\theta\right) \right\rangle$$
$$= \left\langle f_{\xi}\left(a, u\left(a\right), \dot{u}\left(a\right)\right); \frac{d}{ds} \left[\bar{\phi}_{s} \circ u\right] \Big|_{s=0}\left(a\right) \right\rangle.$$

Since $\theta \in (a, b]$ is arbitrary, this proves the special case.

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Now we are going to prove the general case.

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Now we are going to prove the general case. We want to reduce it to the special case we just proved.

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Now we are going to prove the general case. We want to reduce it to the special case we just proved. So we artifically introduce a variable τ

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Now we are going to prove the general case. We want to reduce it to the special case we just proved. So we artifically introduce a variable τ and consider t as a new dependent variable on the same footing as u

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Now we are going to prove the general case. We want to reduce it to the special case we just proved. So we artifically introduce a variable τ and consider t as a new dependent variable on the same footing as u to transform the problem to the previous case, but on \mathbb{R}^{N+1} instead.

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Now we are going to prove the general case. We want to reduce it to the special case we just proved. So we artifically introduce a variable τ and consider t as a new dependent variable on the same footing as u to transform the problem to the previous case, but on \mathbb{R}^{N+1} instead.

For $au \in [a, b]$, we write t = t(au) := au and set the new Lagrangian density

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For $au \in [a, b]$, we write t = t(au) := au and set the new Lagrangian density

$$\bar{f}\left(t\left(\tau\right), \ u\left(t\left(\tau\right)\right), \ \frac{dt}{d\tau}, \ \frac{d}{d\tau}\left[u\left(t\left(\tau\right)\right)\right]\right)$$

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Now we are going to prove the general case. We want to reduce it to the special case we just proved. So we artifically introduce a variable τ and consider t as a new dependent variable on the same footing as u to transform the problem to the previous case, but on \mathbb{R}^{N+1} instead.

For $au \in [a, b]$, we write t = t(au) := au and set the new Lagrangian density

$$\overline{f}\left(t\left(\tau\right), \ u\left(t\left(\tau\right)\right), \ \frac{dt}{d\tau}, \ \frac{d}{d\tau}\left[u\left(t\left(\tau\right)\right)\right]\right)$$
$$:= f\left(t, u\left(t\right), \frac{\frac{d}{d\tau}\left[u\left(t\left(\tau\right)\right)\right]}{\frac{dt}{d\tau}}\right) \frac{dt}{d\tau}$$

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Now we are going to prove the general case. We want to reduce it to the special case we just proved. So we artifically introduce a variable τ and consider t as a new dependent variable on the same footing as u to transform the problem to the previous case, but on \mathbb{R}^{N+1} instead.

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$$:= f\left(t, u\left(t\right), \frac{\frac{d}{d\tau}\left[u\left(t\left(\tau\right)\right)\right]}{\frac{dt}{d\tau}}\right) \frac{dt}{d\tau}$$
$$= f\left(t, u\left(t\right), \dot{u}\left(t\right)\right) \frac{dt}{d\tau}.$$

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Now we are going to prove the general case. We want to reduce it to the special case we just proved. So we artifically introduce a variable τ and consider t as a new dependent variable on the same footing as u to transform the problem to the previous case, but on \mathbb{R}^{N+1} instead.

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$$:= f\left(t, u\left(t\right), \frac{\frac{d}{d\tau}\left[u\left(t\left(\tau\right)\right)\right]}{\frac{dt}{d\tau}}\right) \frac{dt}{d\tau}$$
$$= f\left(t, u\left(t\right), \dot{u}\left(t\right)\right) \frac{dt}{d\tau}.$$

So the Lagrangian

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Now we are going to prove the general case. We want to reduce it to the special case we just proved. So we artifically introduce a variable τ and consider t as a new dependent variable on the same footing as u to transform the problem to the previous case, but on \mathbb{R}^{N+1} instead.

For $au \in [a, b]$, we write t = t(au) := au and set the new Lagrangian density

$$\bar{f}\left(t\left(\tau\right), \ u\left(t\left(\tau\right)\right), \ \frac{dt}{d\tau}, \ \frac{d}{d\tau}\left[u\left(t\left(\tau\right)\right)\right]\right)$$
$$:= f\left(t, u\left(t\right), \frac{\frac{d}{d\tau}\left[u\left(t\left(\tau\right)\right)\right]}{\frac{dt}{d\tau}}\right) \frac{dt}{d\tau}$$
$$= f\left(t, u\left(t\right), \dot{u}\left(t\right)\right) \frac{dt}{d\tau}.$$

So the Lagrangian

$$\bar{I}(t,u) := \int_{\tau_0}^{\tau_1} \bar{f}\left(t\left(\tau\right), \ u\left(t\left(\tau\right)\right), \ \frac{dt}{d\tau}, \ \frac{d}{d\tau}\left[u\left(t\left(\tau\right)\right)\right]\right) \ \mathrm{d}\tau$$

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Now we are going to prove the general case. We want to reduce it to the special case we just proved. So we artifically introduce a variable τ and consider t as a new dependent variable on the same footing as u to transform the problem to the previous case, but on \mathbb{R}^{N+1} instead.

For $au \in [a, b]$, we write t = t(au) := au and set the new Lagrangian density

$$\bar{f}\left(t\left(\tau\right), \ u\left(t\left(\tau\right)\right), \ \frac{dt}{d\tau}, \ \frac{d}{d\tau}\left[u\left(t\left(\tau\right)\right)\right]\right)$$
$$:= f\left(t, u\left(t\right), \frac{\frac{d}{d\tau}\left[u\left(t\left(\tau\right)\right)\right]}{\frac{dt}{d\tau}}\right) \frac{dt}{d\tau}$$
$$= f\left(t, u\left(t\right), \dot{u}\left(t\right)\right) \frac{dt}{d\tau}.$$

So the Lagrangian

$$\bar{I}(t,u) := \int_{\tau_0}^{\tau_1} \bar{f}\left(t(\tau), u(t(\tau)), \frac{dt}{d\tau}, \frac{d}{d\tau}[u(t(\tau))]\right) d\tau$$

does not depend explicitly on τ anymore.

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Thus we can apply the previous result to \overline{I} and deduce that

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$$\left\langle \bar{f}_{\bar{\xi}}\left(t\left(\tau\right), \ u\left(t\left(\tau\right)\right), \ \frac{dt}{d\tau}, \ \frac{d}{d\tau}\left[u\left(t\left(\tau\right)\right)\right]\right); \ \frac{d}{ds}\left[\bar{\phi}_{s}\circ u\right]\Big|_{s=0}\left(t\left(\tau\right)\right)\right\rangle$$

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Thus we can apply the previous result to \overline{I} and deduce that

$$\left\langle \bar{f}_{\bar{\xi}} \left(t\left(\tau\right), \ u\left(t\left(\tau\right)\right), \ \frac{dt}{d\tau}, \ \frac{d}{d\tau} \left[u\left(t\left(\tau\right)\right) \right] \right); \ \frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \bigg|_{s=0} \left(t\left(\tau\right) \right) \right\rangle + \bar{f}_{\xi^{0}} \left(t\left(\tau\right), \ u\left(t\left(\tau\right)\right), \ \frac{dt}{d\tau}, \ \frac{d}{d\tau} \left[u\left(t\left(\tau\right)\right) \right] \right) \ \frac{d}{ds} \left[\phi_{s}^{0} \circ t \right] \bigg|_{s=0} \left(\tau \right) \right)$$

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econd Variation

Thus we can apply the previous result to \overline{I} and deduce that

$$\left\langle \bar{f}_{\bar{\xi}} \left(t\left(\tau\right), \ u\left(t\left(\tau\right)\right), \ \frac{dt}{d\tau}, \ \frac{d}{d\tau} \left[u\left(t\left(\tau\right)\right) \right] \right); \ \frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \Big|_{s=0} \left(t\left(\tau\right) \right) \right\rangle \right. \\ \left. \left. + \bar{f}_{\xi^{0}} \left(t\left(\tau\right), \ u\left(t\left(\tau\right)\right), \ \frac{dt}{d\tau}, \ \frac{d}{d\tau} \left[u\left(t\left(\tau\right)\right) \right] \right) \ \frac{d}{ds} \left[\phi_{s}^{0} \circ t \right] \right|_{s=0} \left(\tau \right) \right\}$$

is a first integral,

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Thus we can apply the previous result to \overline{I} and deduce that

$$\left\langle \bar{f}_{\bar{\xi}} \left(t\left(\tau\right), \ u\left(t\left(\tau\right)\right), \ \frac{dt}{d\tau}, \ \frac{d}{d\tau} \left[u\left(t\left(\tau\right)\right) \right] \right); \ \frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \Big|_{s=0} \left(t\left(\tau\right) \right) \right\rangle$$

$$+ \bar{f}_{\xi^{0}} \left(t\left(\tau\right), \ u\left(t\left(\tau\right)\right), \ \frac{dt}{d\tau}, \ \frac{d}{d\tau} \left[u\left(t\left(\tau\right)\right) \right] \right) \ \frac{d}{ds} \left[\phi_{s}^{0} \circ t \right] \Big|_{s=0} \left(\tau\right) \right\}$$

is a first integral, where ξ^0 is the $dt/d\tau$ variable and $\bar{\xi}$ stands for the $d\left[u\left(t\left(\tau\right)\right)\right]/d\tau$ variable.

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Thus we can apply the previous result to \overline{I} and deduce that

$$\left\langle \bar{f}_{\xi} \left(t\left(\tau\right), \ u\left(t\left(\tau\right)\right), \ \frac{dt}{d\tau}, \ \frac{d}{d\tau} \left[u\left(t\left(\tau\right)\right) \right] \right); \ \frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \bigg|_{s=0} \left(t\left(\tau\right) \right) \right\rangle$$

$$+ \bar{f}_{\xi^{0}} \left(t\left(\tau\right), \ u\left(t\left(\tau\right)\right), \ \frac{dt}{d\tau}, \ \frac{d}{d\tau} \left[u\left(t\left(\tau\right)\right) \right] \right) \ \frac{d}{ds} \left[\phi_{s}^{0} \circ t \right] \bigg|_{s=0} \left(\tau \right) \right\rangle$$

is a first integral, where ξ^0 is the $dt/d\tau$ variable and $\bar{\xi}$ stands for the $d\left[u\left(t\left(\tau\right)\right)\right]/d\tau$ variable. Now from the definition of \bar{f} , clearly at s = 0, we have

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$$\left\langle \bar{f}_{\xi} \left(t\left(\tau\right), \ u\left(t\left(\tau\right)\right), \ \frac{dt}{d\tau}, \ \frac{d}{d\tau} \left[u\left(t\left(\tau\right)\right) \right] \right); \ \frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \bigg|_{s=0} \left(t\left(\tau\right) \right) \right\rangle$$

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 $\bar{f}_{\bar{\xi}} = f_{\xi}$

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$$\bar{f}_{\bar{\xi}} = f_{\xi}$$

and we compute

$$ar{f}_{\xi^0} = f - \langle f_{\xi}; \dot{u}
angle$$
.

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. (since $\phi^0_s\left(t\left(au
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ight)= au,$ we have $dt/d au=1$ at $s=0$)

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$$\left\langle \bar{f}_{\xi} \left(t\left(\tau\right), \ u\left(t\left(\tau\right)\right), \ \frac{dt}{d\tau}, \ \frac{d}{d\tau} \left[u\left(t\left(\tau\right)\right) \right] \right); \ \frac{d}{ds} \left[\bar{\phi}_{s} \circ u \right] \Big|_{s=0} \left(t\left(\tau\right) \right) \right\rangle$$

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This proves the theorem.

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Examples of conservation laws

Now we illustrate Noether's theorem by some simple examples.

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Now we illustrate Noether's theorem by some simple examples.

Example 1: Linear momentum conservation Let m > 0 be the mass and $x(t) \in \mathbb{R}^3$ be the position of a point particle. Let the potential energy function $U : \mathbb{R}^3 \to \mathbb{R}$ be independent of the x_3 coordinate, i.e.

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$$U(x)=U(x_1,x_2).$$

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The Lagrangian density, as usual, is

$$f(t,x,\xi)=\frac{1}{2}m\xi^2-U(x).$$

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It is easy to check that the family of transformations

$$h_{s}(x) := x + se_{3} = (x_{1}, x_{2}, x_{3} + s), \qquad s \in \mathbb{R}$$

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$$p_3 := m\dot{x}_3$$

is a first integral. Similarly, if $U \equiv 0$, then $m\dot{x}$ is a first integral.

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Example 2: Angular momentum conservation Once again consider the same Lagrangian density.

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Example 2: Angular momentum conservation Once again consider the same Lagrangian density. But this time let the potential energy function $U : \mathbb{R}^3 \to \mathbb{R}$ satisfy

 $U(x) = U(R_3[\phi]x)$ for all $\phi \in [0, 2\pi]$,

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$$R_3[\phi] = \begin{pmatrix} \cos\phi & -\sin\phi & 0\\ \sin\phi & \cos\phi & 0\\ 0 & 0 & 1 \end{pmatrix}$$

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leaves the Lagrangian invariant. Now one can check that Noether's theorem tells us that the third component of angular momentum

$$M_3 := (u \wedge m\dot{x}) \cdot e_3$$

is a first integral, where \wedge is the vector product (cross product).

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The first of our example can be generalized in a sense. Assume we are given a Hamiltonian which has no explicit dependence,

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H(t, u, v) = H(u, v).

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The first of our example can be generalized in a sense. Assume we are given a Hamiltonian which has no explicit dependence, i.e.

H(t, u, v) = H(u, v).

The variables u and v in the arguments of the Hamiltonian are in a sense **conjugate variables**.

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The variables u and v in the arguments of the Hamiltonian are in a sense **conjugate variables**. For each $1 \le i \le N$, the variable v_i , is called the **conjugate momenta** of the variable u_i (for the same i).

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$$\frac{\partial H}{\partial u_i}=0.$$

Then the variable u_i is called a **cyclic variable**.

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Then the variable u_i is called a **cyclic variable**. See the notes for a precise definition. Now it can be shown that the Noether's theorem implies

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Then the variable u_i is called a **cyclic variable**. See the notes for a precise definition. Now it can be shown that the Noether's theorem implies that the conjugate momenta for a cyclic variable is a first integral.

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Then the variable u_i is called a **cyclic variable**. See the notes for a precise definition. Now it can be shown that the Noether's theorem implies that the conjugate momenta for a cyclic variable is a first integral. More precisely,

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See assignments and notes for more on this.

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The End

Thank you *Questions?*